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Integrability of the vacuum Einstein equation

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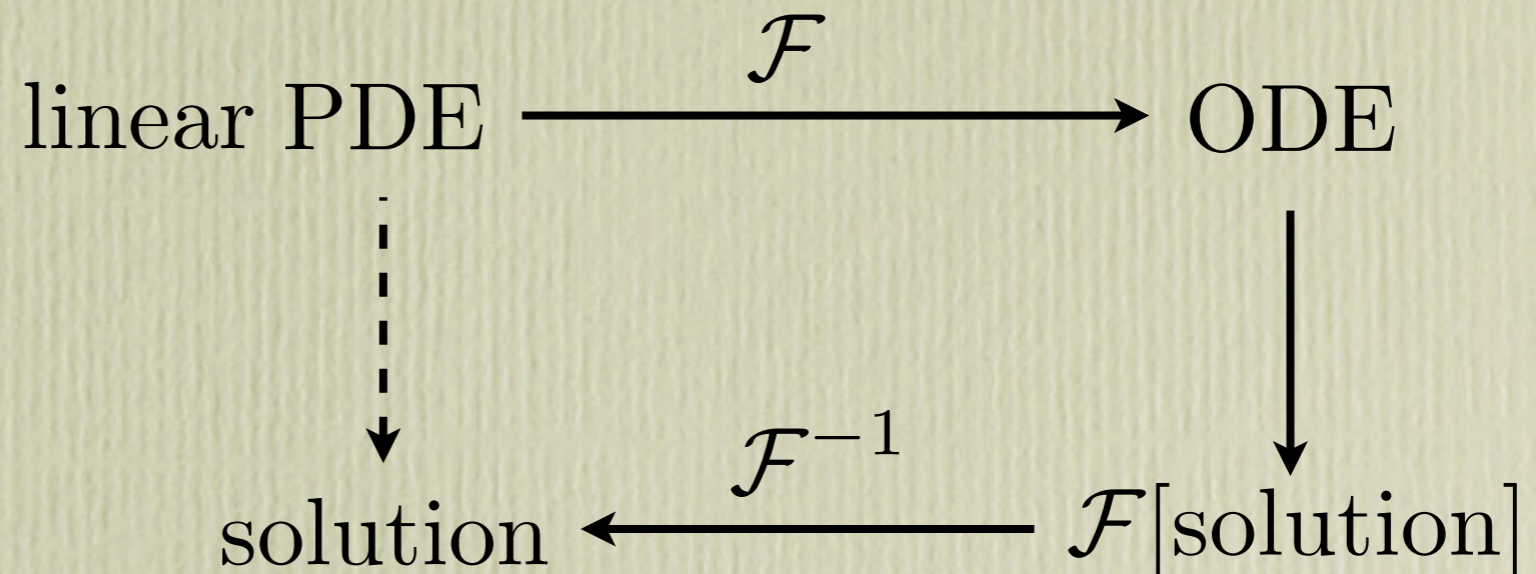
The Inverse Spectral Transform



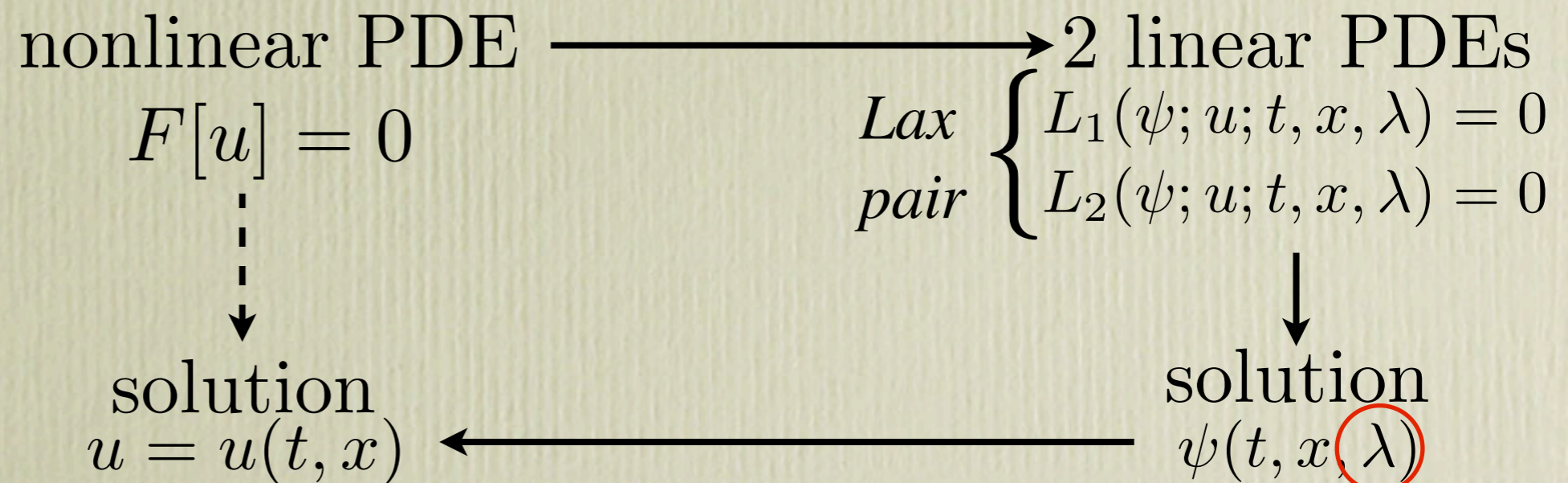
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The **Inverse Scattering Transform** (IST) is one of the most important developments in mathematical physics in the past 40 years. The method is a non-linear analogue, and in some sense generalization, of the [Fourier transform](#), ...

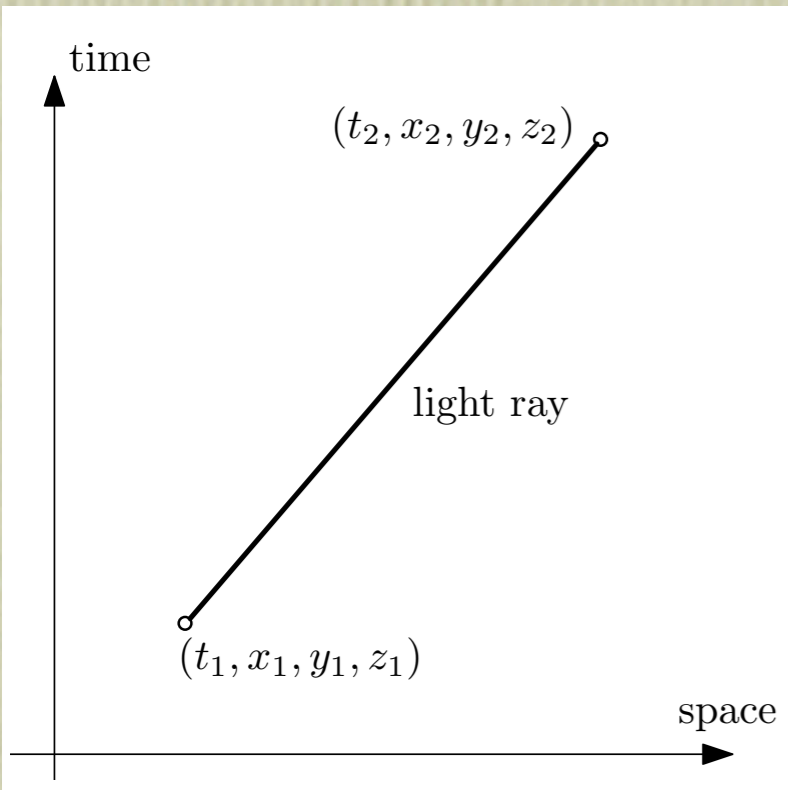
The Fourier Transform



The Inverse Scattering Transform



The theory of relativity



$$\sqrt{(dx)^2 + (dy)^2 + (dz)^2} = dt \iff ds^2 = 0$$

Denote $x^0 = t, x^1 = x, x^2 = y, x^3 = z$

and $\eta_{\mu\nu} = \begin{bmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$

The spacetime interval $ds^2 := g_{\mu\nu} dx^\mu dx^\nu$

The equivalence principle: gravity affects all bodies in the same way, independently of their composition.

In empty space, Spacetime curves according to Einstein's vacuum equation

Ricci Curvature $R_{\mu\nu} = 0$

The Belinski-Zakharov transform

$$ds^2 = f(-dt^2 + dr^2) + g_{ab}dx^a dx^b \quad a, b = 1, 2 \quad x^a = (\phi, z)$$

In matrix form, $g_{\mu\nu} = \begin{bmatrix} -f & 0 & 0 & 0 \\ 0 & g_{11} & g_{12} & 0 \\ 0 & g_{21} & g_{22} & 0 \\ 0 & 0 & 0 & f \end{bmatrix}$ light-cone coordinates
 $t = \zeta - \eta \quad r = \zeta + \eta$

$$R_{\mu\nu} = 0 \quad \underbrace{(\alpha g_{,\zeta} g^{-1})_{,\eta}}_{-A} + \underbrace{(\alpha g_{,\eta} g^{-1})_{,\zeta}}_B = 0 \quad \det(g) = \alpha^2$$

f is easily integrated once g is known

$$\left. \begin{aligned} D_1 &= \partial_\zeta - \frac{2\alpha_{,\zeta}\lambda}{\lambda - \alpha} \partial_\lambda \\ D_2 &= \partial_\eta + \frac{2\alpha_{,\eta}\lambda}{\lambda + \alpha} \partial_\lambda \end{aligned} \right\} \text{Commuting operators} \quad \left. \begin{aligned} D_1\psi &= \frac{A}{\lambda - \alpha} \psi \\ D_2\psi &= \frac{B}{\lambda + \alpha} \psi \end{aligned} \right\} \text{Lax pair}$$

In the limit $\lambda \rightarrow 0$ get $g(\zeta, \eta) = \psi(0, \zeta, \eta)$

The dressing method

1

Take a particular
“background”
solution $g^{(0)}(\zeta, \eta)$

can skip this step

2

$$\begin{aligned} D_1 \psi_0 &= \frac{A_0}{\lambda - \alpha} \psi_0 \\ D_2 \psi_0 &= \frac{B_0}{\lambda + \alpha} \psi_0 \end{aligned} \implies \text{get } \psi_0$$

3

Search for a solution of
the form $\psi = \chi \psi_0$

χ : the dressing matrix

$$\chi = I + \sum_{s=1}^N \frac{R_s}{\lambda - \mu_s}$$

4

Get the N-soliton solution

$$g^{(N)} = \chi(0) g^{(0)}$$

‘Cylindrical’ metrics: $ds^2 = f^{(0)}(-dt^2 + dr^2) + e^{\Lambda^{(0)}}(rd\phi)^2 + e^{-\Lambda^{(0)}} dz^2$

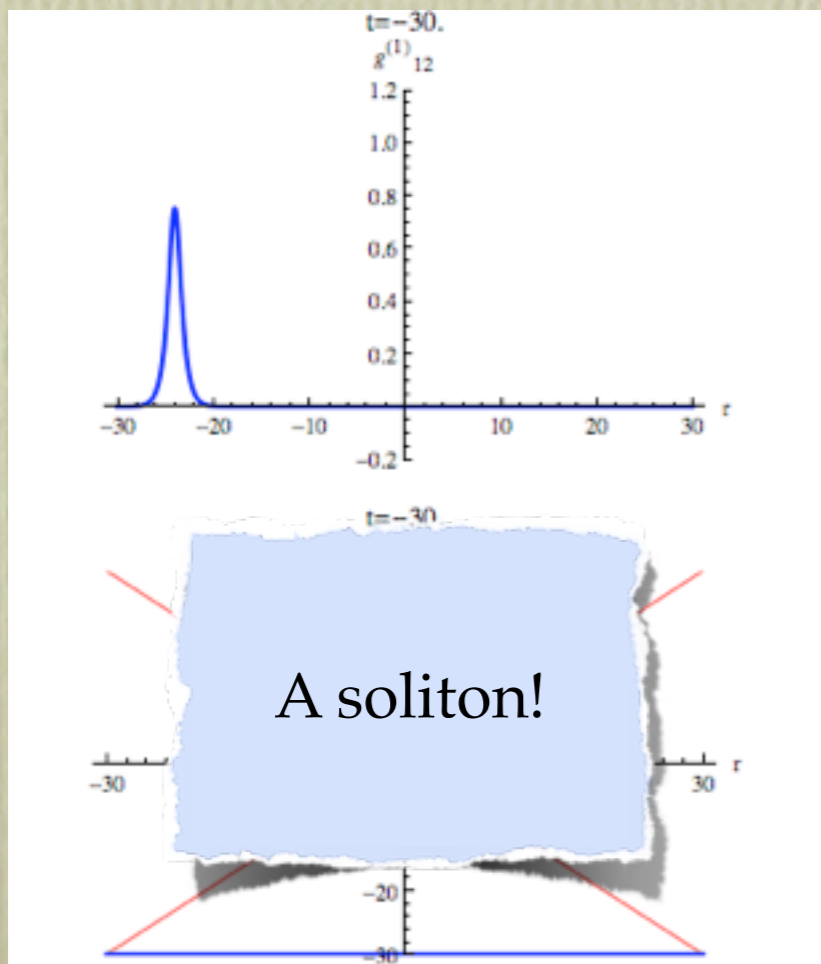
Solitons (one soliton)

Symmetric chiral equation

$$(g, \zeta g^{-1})_{,\eta} + (g, \eta g^{-1})_{,\zeta} = 0$$

$$g^{(0)} = \begin{bmatrix} e^{\Lambda^{(0)}} & 0 \\ 0 & e^{-\Lambda^{(0)}} \end{bmatrix}$$

$$\mu_s = \text{constant}$$

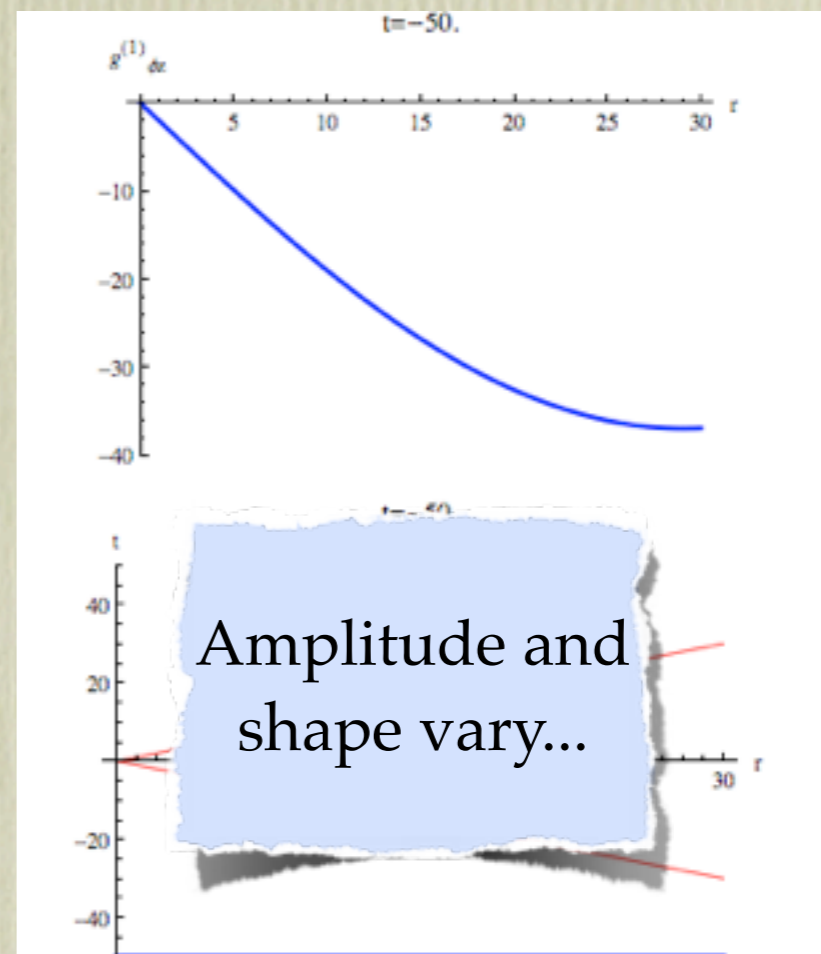


Einstein equations

$$(\alpha g, \zeta g^{-1})_{,\eta} + (\alpha g, \eta g^{-1})_{,\zeta} = 0$$

$$g^{(0)} = \begin{bmatrix} r^2 e^{\Lambda^{(0)}} & 0 \\ 0 & e^{-\Lambda^{(0)}} \end{bmatrix}$$

$$\mu_s = \omega_s - t \pm \sqrt{(\omega_s - t)^2 - r^2}$$



Solitons (two solitons)

For physical reasons, one would like the metric to be asymptotically flat.

The one-soliton solution is never asymptotically flat.

Pole $\mu_s = \omega_s - t \oplus \sqrt{(\omega_s - t)^2 - r^2}$ $+ = \text{soliton}$
 $- = \text{anti-soliton}$

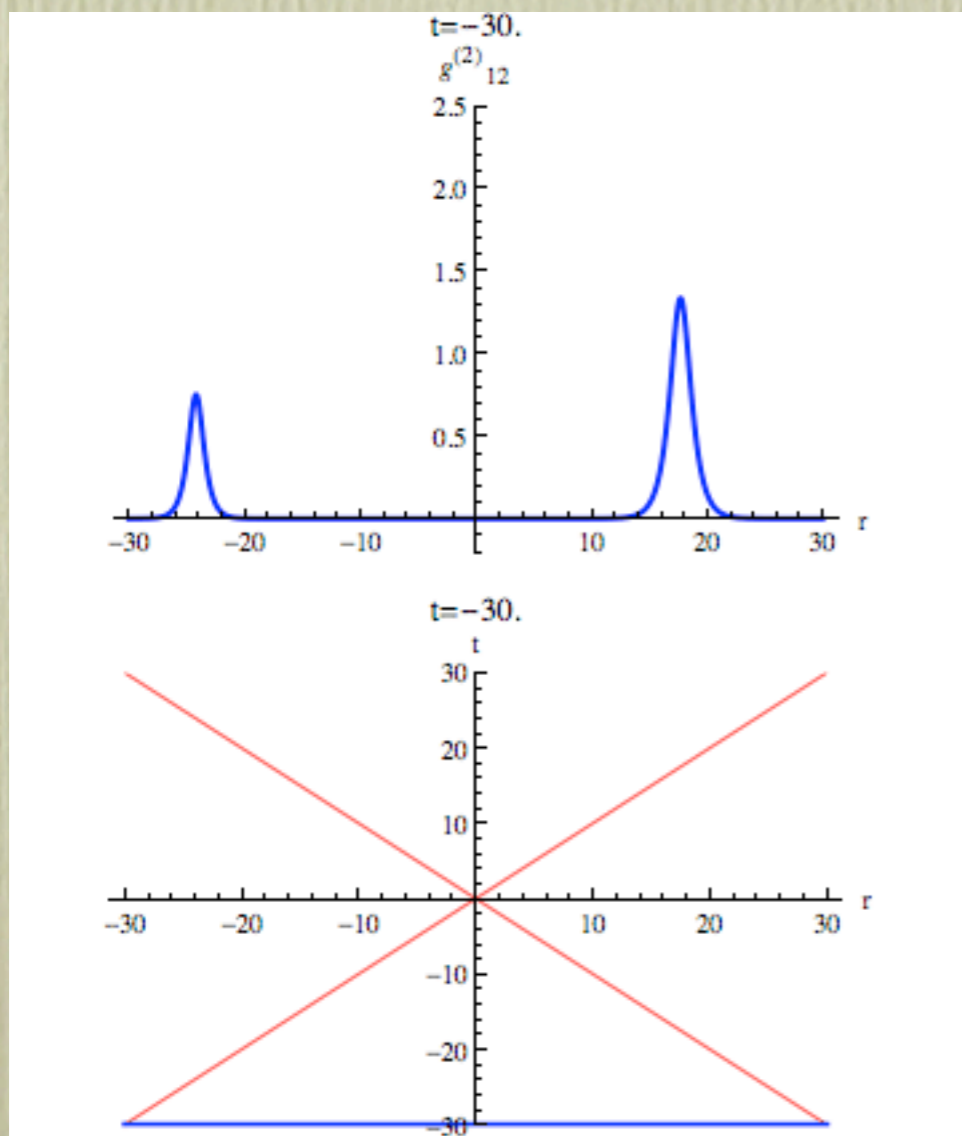
Theorem: If $\Lambda^{(0)} \xrightarrow{|t| \rightarrow \infty} 0$ and $(\# \text{ of solitons}) = (\# \text{ of anti-solitons})$ then the solution is asymptotically flat as $|t| \rightarrow \infty$.

Both S-S or AS-AS are never asymptotically flat.

Solitons (two solitons)

Symmetric chiral equation

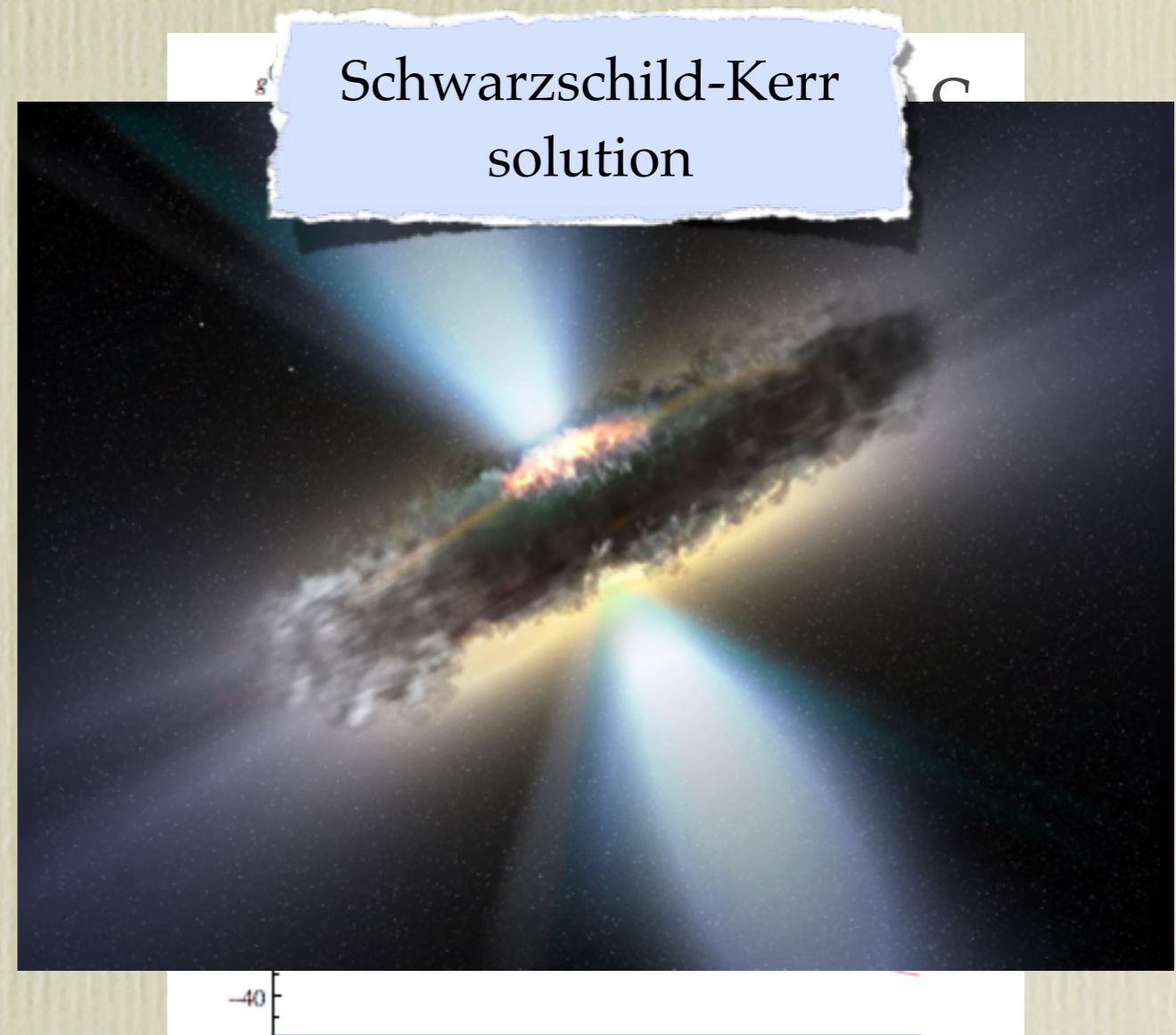
$$g^{(0)} = \begin{bmatrix} e^{\Lambda^{(0)}} & 0 \\ 0 & e^{-\Lambda^{(0)}} \end{bmatrix}$$



Einstein equations

$$g^{(0)} = \begin{bmatrix} r^2 e^{\Lambda^{(0)}} & 0 \\ 0 & e^{-\Lambda^{(0)}} \end{bmatrix}$$

$\Lambda^{(0)} = 0$ (flat)



Thank you!



Supplemental Slides

